# Traveling pulses and branching patterns Benoît Perthame



# WHY



Bacterial colonies. Top S. Serror (CNRS-Paris-Sud). Bottom K. Ben Jacob (TAU)

#### WHY

Lecture 1. Parabolic models and pulse propagation

Lecture 2. The hyperbolic Keller-Segel model and branching pattern

Lecture 2bis. V. Calvez (Keller-Segel and asymmetric pulses)

Lecture 3. Another example of concentration Darwin evolution

### MIMURA's model

$$\begin{cases} \frac{\partial}{\partial t}n(t,x) - d_1\Delta n = r n\left(S - \frac{\mu n}{(n_0 + n)(S_0 + S)}\right),\\\\ \frac{\partial}{\partial t}S(t,x) - d_2\Delta S = -r nS,\\\\ \frac{\partial}{\partial t}f(t,x) = r n \frac{\mu n}{(n_0 + n)(S_0 + S)}\end{cases}$$

The dynamics is driven by the source terms, i.e., by bacterial growth/nutrient consumption.



## Intuitive explanation

• Nutrient is consumed by the active cells and reaches a low level in the colony

- Cells at the front have an advantage for multiplication
- This enhences any perturbbation

This model, as well as other variants are based on the Gray-Scott chemical reaction

$$\frac{\partial}{\partial t}u - d_u \Delta u = u \left( u^{n-1}v - \mu \right),$$
$$\frac{\partial}{\partial t}v - d_v \Delta v = -u^n v,$$
$$\frac{\partial}{\partial t}f(t, x) = \mu u^n.$$

Here n = 1, 2... plays the role of ignititon temperature.

Levin and Kessler model is

$$\frac{\partial}{\partial t}u - d_u \Delta u = u (h(u)v - \mu),$$
  
$$h(u) = 0 \qquad for \qquad u < u_{\text{threshold}}$$

•

This model, as well as other variants are based on the Gray-Scott chemical reaction

$$\begin{split} \frac{\partial}{\partial t}u - d_u \Delta u &= u \Big( u^{n-1}v - \mu \Big), \\ \frac{\partial}{\partial t}v - d_v \Delta v &= -u^n v, \\ \frac{\partial}{\partial t}f(t, x) &= \mu u^n. \end{split}$$

Here n = 1 makes a big difference with n = 2... or Levin and Kessler model or Mimura model.

Traveling pulse for Gray-Scott

$$\begin{cases} -\sigma u' - d_u u'' = u (v - \mu), & u(\pm \infty) = 0, \\ -\sigma v' = -uv, & v(-\infty) = v_-, & v(+\infty) = v_+ \end{cases}$$

**Theorem (PES, BP ongoing)** Let  $(\mu, v_-, v_+)$  be such that

$$v_{-} < \mu < v_{+}, \qquad \mu \ln(v_{-}) - v_{-} = \mu \ln(v_{+}) - v_{+},$$

Then, for all speeds

$$\sigma > \sigma^* := 2\sqrt{v_+ - \mu}$$

there is a unique traveling pulse and v is increasing.

Some references

Golding-Koslovsky-Ben Jacob

Muratov-Osipov, Doelman-Eckhaus-Kaper-Gardner

M. Ward-Kolokolnikov-Wei

Mimura (Masuda)

Elliptic case : Del Pino, Kowalckzyk



Pulse splitting in GS system (from DoelmanEckhaus-Kaper, SIAP



Pulse splitting in 2D GS system (from M. Ward)



Pulse splitting in 2D GS system (from M. Ward)

New experiments are done on rich media.

Are there models based on other ingredients that achieve this type of patterns?



Experiments by B. Holland and S. Serror, CNRS Paris-Sud

The shortcoming of Keller-Segel system

$$\begin{cases} \frac{\partial}{\partial t}n(t,x) + \operatorname{div}[\chi n \nabla c] = \Delta n, & x \in \mathbb{R}^d, \ t \ge 0, \\ -\Delta c + c = n, \\ n(t,x) = n^0(x). \end{cases}$$

The shortcoming of Keller-Segel system

$$\begin{bmatrix} \frac{\partial}{\partial t}n(t,x) + \operatorname{div}[\chi n\nabla c] = \Delta n + n(1-n), & x \in \mathbb{R}^d, \ t \ge 0, \\ -\Delta c + c = n, \\ n(t,x) = n^0(x). \end{bmatrix}$$

## Theorem (G. Nadin, BP, L. Ryzhik)

For  $\chi$  small enough, there is a traveling wave.

For  $\chi$  large enough the K.-S. equation is unstable in the sense of Turing.

The hyperbolic Keller-Segel system (Dolak, Schmeiser)

$$\begin{cases} \frac{\partial}{\partial t}n(t,x) + \operatorname{div}[n(1-n)\nabla c] = 0, & x \in \mathbb{R}^d, \ t \ge 0, \\ -\Delta c + c = n, \\ n(t,x) = n^0(x), & 0 \le n^0(x) \le 1, \quad n^0 \in L^1(\mathbb{R}^d). \end{cases}$$

Interpretation

- n(t, x) = bacterial density ,
- c(t, x) = chemical signalling (chemoattraction),
- n(1-n) represents something like quorum sensing,
- random motion of bacterials is neglected

Related to the Keller-Segel model but no point concentrations



By V. Calvez, B. Desjardins, H. Khonsari on multiple sclerosis

$$\frac{\partial}{\partial t}n(t,x) + \operatorname{div}\left[n(1-n)\nabla c - n\nabla S\right] = 0, \qquad x \in \mathbb{R}^d, \ t \ge 0,$$
$$-d_c\Delta c + c = \alpha_c n,$$
$$\frac{\partial}{\partial t}f(t,x) - d_f\Delta f = \alpha_f n + f(1-f)$$
$$\frac{\partial}{\partial t}S(t,x) - d_S\Delta S + S = \alpha_S \left(n_{mother\ colony} + f + n\right).$$

- n = swarmer cells,
- f = follower (supporter) cells,
- c = short range 'attractant'
- S = long range signal (surfactant?)

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Numerical instabilities can be observed on reduced systems

$$\begin{cases} \frac{\partial}{\partial t}n(t,x) + \operatorname{div}\left[n(1-n)\nabla c - n\nabla S\right] = 0, & x \in \mathbb{R}^d, \ t \ge 0, \\ -d_c\Delta c + c = \alpha_c n, \\ \frac{\partial}{\partial t}S(t,x) - d_S\Delta S + \tau_S S = \alpha_S n. \end{cases}$$

And check for traveling pulses

$$\begin{cases} -\sigma n' + \left[n(1-n)c' - nS'\right]' = 0, & x \in \mathbb{R}, \\ -d_c c'' + c = \alpha_c n, \\ -\sigma S' - d_S S'' + \tau_S S = \alpha_S n. \end{cases}$$

$$\begin{cases} -\sigma n' + \left[n(1-n)c' - nS'\right]' = 0, & x \in \mathbb{R}, \quad n(\pm \infty) = 0, \\ -d_c c'' + c = \alpha_c n, \\ -\sigma S' - d_S S'' + \tau_S S = \alpha_S n. \end{cases}$$

$$-\sigma n + n(1-n)c' - nS' = 0, \qquad x \in \mathbb{R},$$
  
 $-d_c c'' + c = \alpha_c n,$   
 $-\sigma S' - d_S S'' + \tau_S S = \alpha_S n.$ 

$$\begin{cases} \bullet \qquad -\sigma + (1-n)c' - S' = 0, \qquad x \in [0, L], \\ \bullet \qquad \qquad n \equiv 0, \text{ for } x \notin [0, L], \\ -d_c c'' + c = \alpha_c n, \\ -\sigma S' - d_S S'' + \tau_S S = \alpha_S n. \end{cases}$$

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• 
$$x \in [0, L],$$
  
•  $n \equiv 0, \text{ for } x \notin [0, L],$   
 $-d_c c'' + c = \alpha_c n,$   
 $-\sigma S' - d_S S'' + \tau_S S = \alpha_S n.$ 

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$$\begin{array}{l} \bullet & -\sigma + (1-n)c' - S' = 0, \qquad x \in [0, L], \\ \bullet & n \equiv 0, \text{ for } x \notin [0, L], \\ -d_c c'' + c = \alpha_c n, \\ -\sigma S' - d_S S'' + \tau_S S = \alpha_S n. \end{array}$$

Special case 1. Steady states,

$$\begin{array}{l} \bullet & (1-n)c'-S'=0, & x\in[0,L], \\ \bullet & n\equiv 0, \ \text{for} \ x\notin[0,L], \\ -d_cc''+c &= \alpha_c n, \\ -d_SS''+S=\alpha_S n. \end{array}$$

#### Theorem For

- L small
- or  $|d_c d_S| + |\alpha_c \alpha_S|$  small

there is a unique solution  $n \in C(0, L)$ .

Special case 1. Steady states,





Special case 2.  $d_S = 0$ ,  $\tau_S = 0$ ,

$$\begin{cases} \bullet & -\sigma + (1-n)c' - S' = 0, \qquad x \in [0, L], \\ \bullet & n \equiv 0, \quad \text{for } x \notin [0, L], \\ -d_c c'' + c = \alpha_c n, \\ -\sigma S' = \alpha_S n. \end{cases}$$

$$n = \left\{ egin{array}{ll} 1, & 0 \leq x \leq L, \\ 0, & ext{otherwise} \end{array} 
ight.$$

 $-\sigma S' = \alpha_S n, \ \forall x, \qquad \sigma = -S' \text{ for } x \in [0, L],$  $\implies \sigma = \sqrt{\alpha_S}, \quad \sigma = -S' \quad x \in [0, L], \qquad S' = 0 \text{ for } x \notin [0, L],$ 

Special case 2.  $d_S = 0$ ,  $\tau_S = 0$  (stability)

Theorem These waves are stable if and only if

$$c'(0) > \sqrt{\alpha}, \qquad \qquad c'(L) < 0.$$

See the problem as an hyperbolic system (as T. LI, Z. WANG).

$$\begin{cases} \frac{\partial}{\partial t}n + [n(1-n)c' - nS_x]_x = 0, & x \in \mathbb{R}, \\ \frac{\partial}{\partial t}S = \alpha_S n. & \\ \begin{cases} \frac{\partial}{\partial t}n + [n(1-n)c' - nv]_x = 0, \\ \frac{\partial}{\partial t}v - \alpha_S n_x = 0. & v := S_x \end{cases}$$

And check the Lax entropy condition.

Special case 2.  $d_S = 0$ ,  $\tau_S = 0$  (structural stability)

Theorem Still when  $c'(0) > \sqrt{\alpha}$ , c'(L) < 0, these waves are stable for  $d_S$  small.



#### Special case 3. $L \approx 0$

$$\begin{cases} \bullet \quad -\sigma + (1-n)c' - S' = 0, \qquad x \in [0, L], \\ \bullet \qquad n \equiv 0, \quad \text{for } x \notin [0, L], \\ -d_c c'' + c = \alpha_c n, \\ -\sigma S' - d_S S'' = \alpha_S n. \end{cases}$$
$$n = 1 + \frac{\sigma + S'}{c'}$$

## Difficulty

- $c'(x_0) = 0$  at a single point
- Choose  $\sigma = S'(x_0)$

Special case 3.  $L \approx 0$ 

$$\begin{cases} n = 1 + \frac{\sigma + S'}{c'} \\ -d_c c'' + c = \alpha_c n, \\ -\sigma S' - d_S S'' = \alpha_S n. \end{cases}$$

**Difficulty** •  $c'(x_0) = 0$  at a single point • Choose  $\sigma = S'(x_0)$ .

Theorem For  $L \approx 0$  there is a unique solution  $n \in C(0, L)$ 

- c is convex in [0, L]
- S is decreasing in [0, L]
- Find a fixed point  $n \mapsto 1 + \frac{\sigma + S'}{c'}$

## Special case 3. $L \approx 0$



## Conclusion

$$\begin{split} \frac{\partial}{\partial t}n(t,x) + \operatorname{div} \Big[n(1-n)\nabla c - n\nabla S\Big] &= 0, \qquad x \in \mathbb{R}^d, \ t \ge 0, \\ -d_c\Delta c + c &= \alpha_c n, \\ \frac{\partial}{\partial t}f(t,x) - d_f\Delta f &= \alpha_f n + f(1-f) \\ \frac{\partial}{\partial t}S(t,x) - d_S\Delta S + S &= \alpha_S \Big(n_{mother\ colony} + f + n\Big). \end{split}$$

This system creates branching patterns.

A reduced hyperbolic Keller-Segel system explains how instabilities can occur on tr solutaveling pulse solutions for n.

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- A.-L. Dalibard, V. Calvez
- C. Schmeiser, M. Tang, N. Vauchelet
- A. Daerr
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